TABLE II. Data for the magnetic B81 to B31 transition in MnAs.

Pressure (kbar):	0.001	1.22	2.33	3.73	4.23	4.5	4.5
$B8_1 \rightarrow B31T_e$ (°K):	317	298	278	243	223	201	77
Temperature (°K):	307	298	283	245	240	201	138
$B31 \rightarrow B8_1 P_e$ (kbar):	0.001	0.7	0.9	2.0	1.7	1.4	0.001

Kelvin. Differentiation gives  $dT_e/dP = -13.8^{\circ}C/kbar$ at 1 atm. This is in moderate agreement with the results of Rodbell and Wilson<sup>8</sup> and of Samara and Giardini,6 who report -12 and  $-12.3^{\circ}C/kbar$ . (DeBlois and Rodbell<sup>3</sup> give -24.1°C/kbar for the  $B31 \rightarrow B8_1$  transition, which is also in reasonable agreement with Table II and Fig. 3.) However, these authors imply that the slope is constant. Our experiments were done on polycrystalline samples, whereas they used single crystals and were confined to P < 2.5 kbar. (Also see note added in proof at end of paper.)

The pressure dependence of  $T_e$  is displayed in Fig. 3, where the smooth curve for increasing pressure in the interval  $190 \le T \le 317^{\circ}$ K is Eq. (2). The data clearly demonstrate a critical pressure  $P_c \approx 4.6$  kbar above which the B31 phase is indeed stable down to lowest temperatures. In addition, the pressure hysteresis increases with decreasing temperature, so that pressures P>2 kbar applied at room temperature are sufficient to stabilize the B31 phase to lowest temperatures. In fact, it is possible to obtain the B31 phase at atmospheric pressure below 138°K if the pressure is reduced from P > 4.6 kbar at a  $T < 138^{\circ}$ K.

The transition at  $T = 138^{\circ}$ K was obtained as follows: A specimen of MnAs was placed under 5-kbar pressure in the high-pressure gas apparatus, and the apparatus was cooled to 77°K. Pressure was then released and the specimen transferred to the cold stage  $(4.2^{\circ}K)$  of a vibrating-coil magnetometer. At no time was the speci-

320

**B31** 

B8, → B31

33 T.2,

is in degrees

(2)

FIG. 2. Electrical istance versus ssure at 25°C of lycrystalline nAs.

51).



'As quoted by Ref. 3.

120°K.

men allowed to warm up more than a few degrees above 77°K. The magnetic properties were measured from 4.2°K to room temperature. Below 138°K there was a small susceptibility that decreased with increasing temperature. An abrupt transition occurred at 138°K. the magnetization increasing by over a factor of 50 to the magnetization value of the  $B8_1$  phase.

Although the low-temperature B31 phase appeared to resemble the metamagnetic phase of low-temperature MnP and MnAs<sub>0.9</sub>P<sub>0.1</sub>, our preliminary measurements on a polycrystalline sample do not allow characterization of the magnetism of this phase. This point is significant because Rodbell and Bean<sup>9</sup> have anticipated a ferromagnetic → antiferromagnetic phase change with increasing pressure at low temperatures. A pressure bomb for further magnetic studies is under construction. Meanwhile, monitoring of the resistance at room temperature up to 12 kbar has shown only one first-order phase change. The Rodbell-Bean  $P^*$ -T phase diagram contains  $P^* = P - \alpha_l T/K$ , where  $\alpha_l$  is the thermal expansion coefficient and K is the compressibility. Their diagram would reflect the  $P^*$ -T curve through the  $P^* = [4.6 + (\alpha_l T/K)]$  kbar line to predict a roomtemperature paramagnetic 
→ antiferromagnetic phase change near  $[8+(\alpha_l T/K)]$  kbar=11.8 kbar. If such a transition exists below 12 kbar, it is not first-order. Further, reduction in temperature under P=3 kbar from the paramagnetic B31 phase at room temperature and 3 kbar gave no anomaly in the resistance down to 90°K. Since the high-pressure phase is magnetically ordered at 90°K, this means that the paramagnetic ₹ magnetic-order transition in the high-pressure phase is not first-order. (The magnetic-order transition temperature for the high-pressure phase is not indicated in Fig. 3 since it is not clearly defined by a resistivity anomaly.) Therefore, it may be assumed that the highpressure phase retains the B31 structure in the areas so designated in Fig. 3 and that there is no crystallographic phase change associated with magnetic ordering in the high-pressure phase.

## **III. DISCUSSION**

## A. Significant Observations

These results, together with the earlier study<sup>5</sup> of  $MnAs_{1-x}P_x$  and the data of Table I, establish the

157

The Institute of

 $(\partial T_e/\partial P)_H > 0;$ 

This is con-

ie electrical

e was moni-

anner. Two

cylinder ap-

transmitting

res, a high-

1 gas as the

atures were

ssel in a con-

peraty run

btaine... The

sure at which

d. This point

n the interval

asing pressure

157

391

<sup>9</sup> D. S. Rodbell and C. P. Bean, J. Appl. Phys. Suppl. 33, 1037 (1962).